

Many-body localization and Spontaneous Symmetry Breaking in the Massless Schwinger Model

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Outline

- 1 Introduction and Motivation
- 2 Lattice models
- 3 Numerical Results
 - Symmetry breaking in the symmetric regularization
 - Ergodicity breaking in the low-energy Hilbert space
 - Localization at low energies
- 4 Concluding remarks
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Some history



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- Motivated by arguments published by Nandkishore and Sondhi (PRX 2017) [1].

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- Goal: investigate lattice formulation (where there is no analytic solution via bosonization), where commensurate effects are also present [7].

How to regularize?

The continuum Hamiltonian formulation in the temporal gauge [8]

$$H = \int dx \{ \bar{\Psi} [i\gamma^1 \partial_1 + e\gamma^1 A_1 + \mu(x)] \Psi + \frac{1}{2} E^2 \} \quad (3)$$

$$\rightarrow H = \frac{-i}{2a} \sum_n (\phi_n^\dagger e^{i\theta_n} \phi_{n+1} - h.c.) + \frac{ag^2}{2} \sum_n L_n^2 \quad (4)$$

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- **Kogut-Susskind staggered (SB) potential (CR)** \iff (2CCP)

$$L_n - L_{n-1} = \phi_n^\dagger \phi_n - \frac{1}{2} [1 - (-1)^n] \quad \langle \phi_n^\dagger \phi_n \rangle = n \pmod{2} \quad (5)$$

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- **Translationally invariant potential (SR)** \iff (1CCP)

$$L_n - L_{n-1} = \phi_n^\dagger \phi_n - 1/2 \quad \langle \phi_n^\dagger \phi_n \rangle = 0 \quad (6)$$

- Jordan-Wigner-ize $\phi_n \rightarrow \prod_{m < n} (i\sigma_m^z) \sigma_n^-$, $\phi_n^\dagger \phi_n \rightarrow \frac{1}{2} (\sigma^z + 1)$

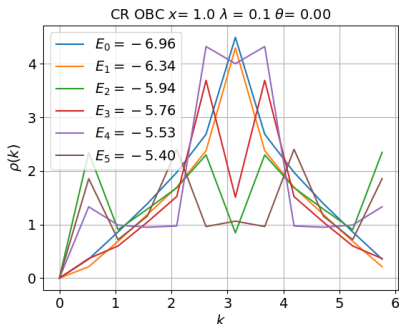
Lattice Hamiltonian

| | Open boundary conditions | Periodic boundary conditions |
|------------------|---|--|
| H_{hop} | $\sum_{n=0}^{N-2} (\sigma_n^+ \sigma_{n+1}^- + \sigma_n^- \sigma_{n+1}^+)$ | $\sum_{n=0}^{N-1} (\sigma_n^+ \sigma_{n+1}^- + \sigma_n^- \sigma_{n+1}^+)$ |
| H_{int} | $\sum_{n=0}^{N-2} (E_n)^2$ | $-\sum_{j<i} \rho(i)\rho(j) \sin\left[\frac{\pi}{N}(i-j)\right]$ |
| Conventional | $E_n = \frac{1}{4} \sum_{k=0}^{N-1} \text{sgn}\left(n + \frac{1}{2} - k\right) \sigma_k^z + \frac{1}{4} [(-1)^n + 1]$ | $\rho(n) = \frac{1}{2} [\sigma_n^z + (-1)^n]$ |
| Symmetric | $E_n = \frac{1}{4} \sum_{k=0}^{N-1} \text{sgn}\left(n + \frac{1}{2} - k\right) \sigma_k^z$ | $\rho(n) = \frac{1}{2} \sigma_n^z$ |

$$H = xH_{\text{hop}} + \lambda H_{\text{int}} + \frac{1}{2} \sum_{n=0}^{N-1} V(n) \sigma_n^z \quad V(n) \in [-\theta, \theta] \quad (7)$$

$$\left(\sum_n \sigma_n^z \right) |j\rangle = 0 \quad (8)$$

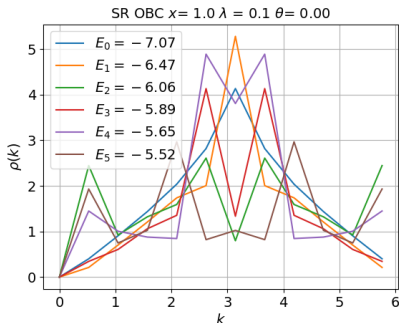
Fourier-transformed density-density correlations ($N = 12$)



$$\rho_n(k) = \sum_{r=0}^{N-1} \langle \sigma_0^z \sigma_r^z \rangle_n e^{ikr} \quad (9)$$

States calculated by E.D. for $N = 12$, Lanczos' algorithm for higher N .

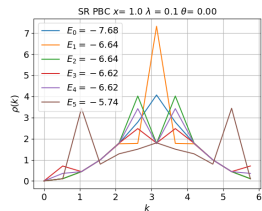
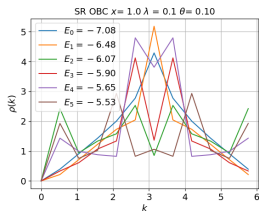
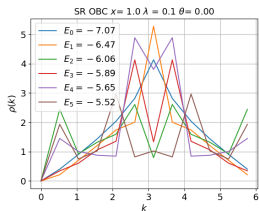
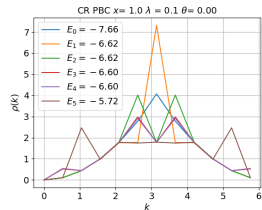
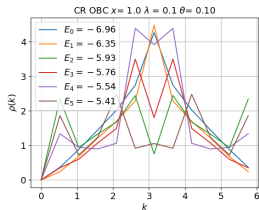
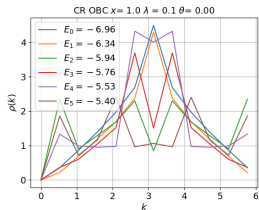
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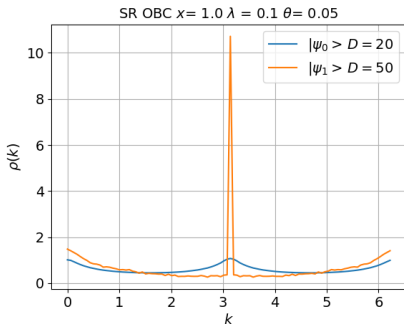
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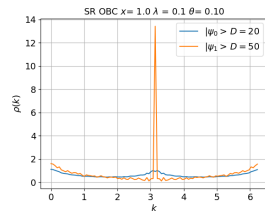
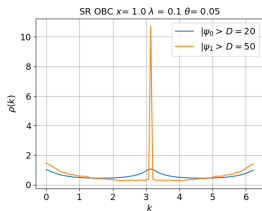
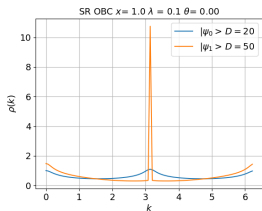
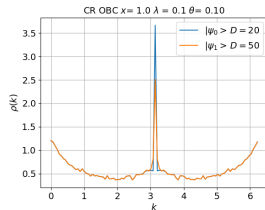
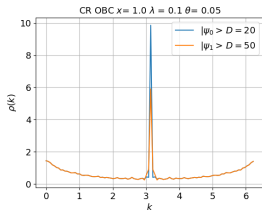
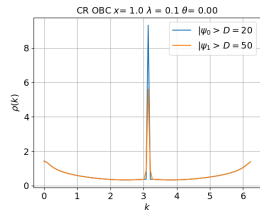
Fourier-transformed density-density correlations ($N = 100$)



$$\rho_n(k) = \sum_{r=0}^{N-1} \langle \sigma_0^z \sigma_r^z \rangle_n e^{ikr} \quad (11)$$

States calculated using iterative ground state search in DMRG. Excited states calculated via penalty, as described in [8].

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The classical 1CCP develops long-range crystalline order at any temperature as well.

What is the level statistics ratio?

$$\tilde{r}_i = \frac{\min(E_{i+2} - E_{i+1}, E_{i+1} - E_i)}{\max(E_{i+2} - E_{i+1}, E_{i+1} - E_i)} \quad (12)$$

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Spectra of time reversal-invariant systems whose classical analogues are K systems show the same fluctuation properties as predicted by GOE.

Level statistics

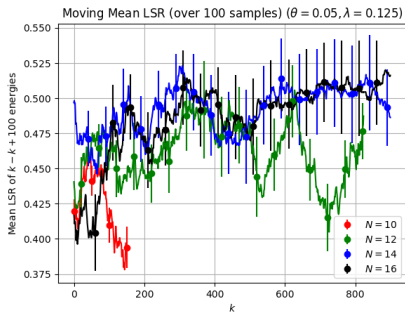
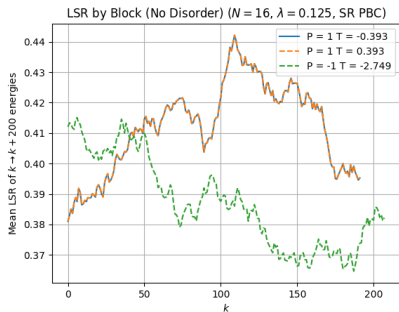
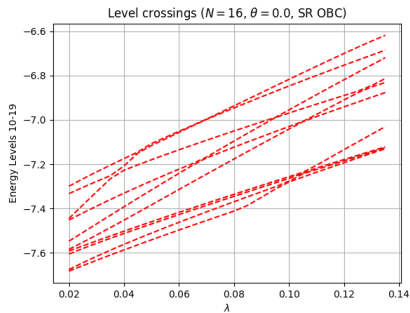
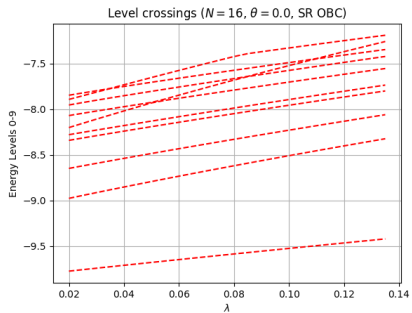
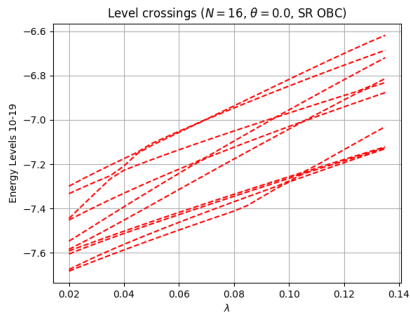
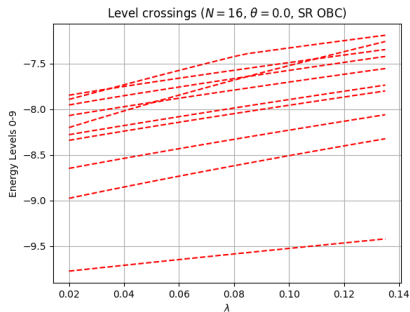


Figure: Window size is 100 energies. Note $\theta = 0$ statistics are far from ergodic. Also, note ergodicity breaking at low energies when we turn on disorder. We speculate this may be because of proximity to an integrable point (the clean continuum Schwinger model). Level crossings in the clean model support this hypothesis.

Level crossings



Level crossings



Since $\xi \approx 20$ according to [1], we must probe larger system sizes.

Density imbalance (CR)

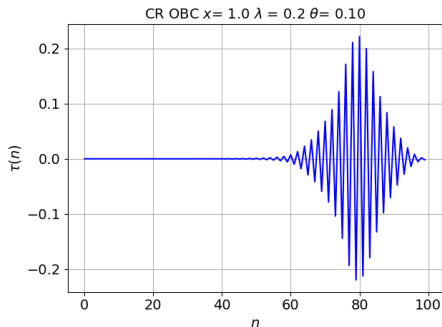
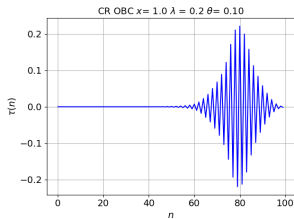
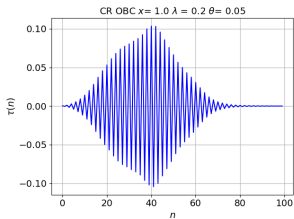
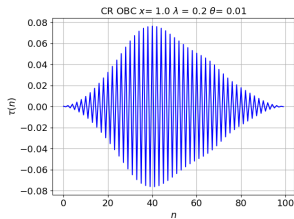
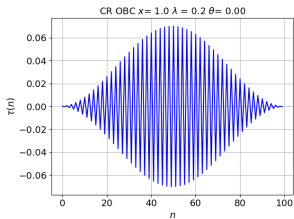


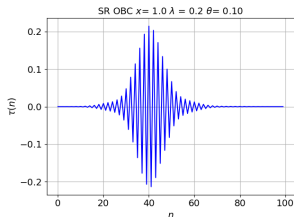
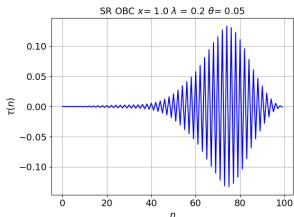
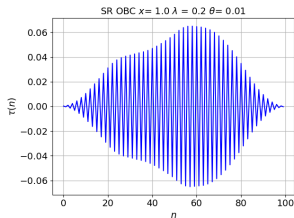
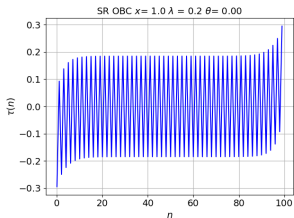
Figure: Localization length shrinks as disorder strength increases. The elementary excitations are charge neutral (dipoles). Parameter regime: $x \gg \theta \sim \lambda$.

$$\tau(n) = \langle \psi_1 | \sigma_n^z | \psi_1 \rangle - \langle \psi_0 | \sigma_n^z | \psi_0 \rangle \quad (15)$$

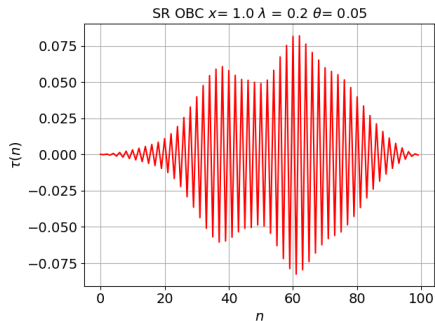
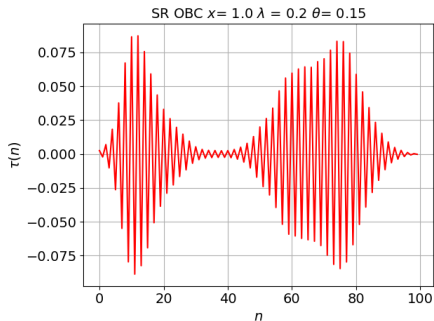
Density imbalance (CR)



Density imbalance (SR)



Bilocalized states



Localization length scaling

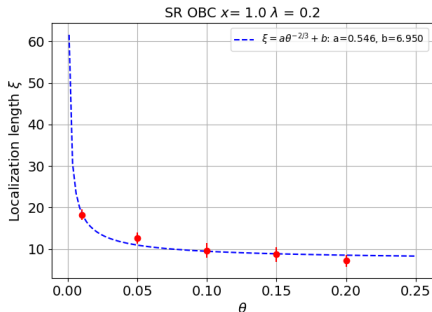
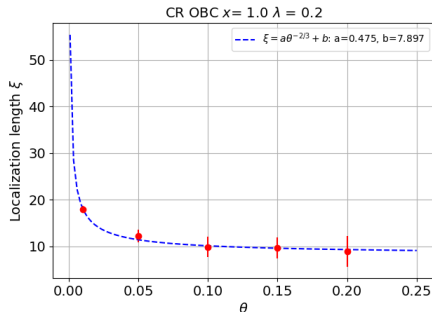


Figure: Comparisons with the predicted localization length of [1].

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- More questions:
 - What happens in other parameter regimes?
 - What happens at higher energies?

-  R. M. Nandkishore and S. L. Sondhi, “Many-body localization with long-range interactions,” *Phys. Rev. X*, vol. 7, p. 041021, Oct 2017.
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